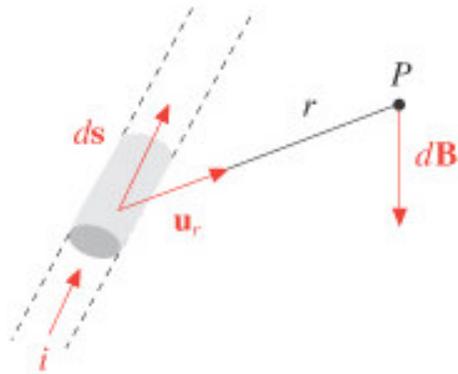


Prima legge di Laplace

$$d\vec{B} = k_m i \frac{d\vec{s} \times \vec{u}_r}{r^2}$$

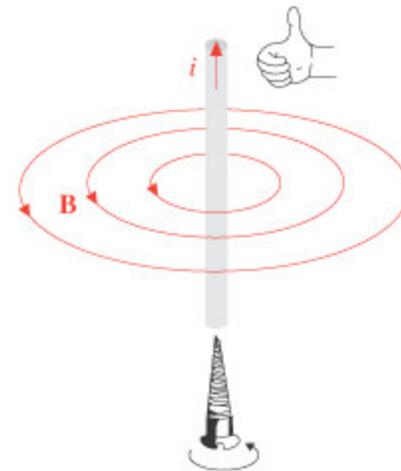


$$k_m = 10^{-7} \frac{Tm}{A} = 10^{-7} \frac{H}{m} = \frac{\mu_0}{4\pi}$$

$$\mu_0 \approx 1.26 \cdot 10^{-6} \frac{H}{m}$$

Prima legge di Laplace

$$d\vec{B} = \frac{\mu_0}{4\pi} i \frac{d\vec{s} \times \vec{u}_r}{r^2}$$



Legge di Ampere-Laplace

$$\vec{B} = \frac{\mu_0}{4\pi} i \oint \frac{d\vec{s} \times \vec{u}_r}{r^2}$$

Campo magnetico di una carica in moto

$$d\vec{B} = \frac{\mu_0}{4\pi} i \frac{d\vec{s} \times \vec{u}_r}{r^2} \xrightarrow{j=nq\vec{v} \quad i=j\Sigma} = \frac{\mu_0}{4\pi} \frac{q\vec{v} \times \vec{u}_r}{r^2} nd\tau$$

Il campo magnetico prodotto da una singola
carica

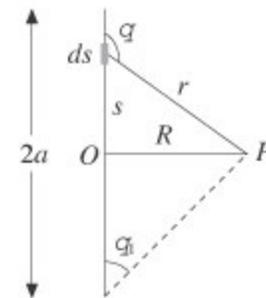
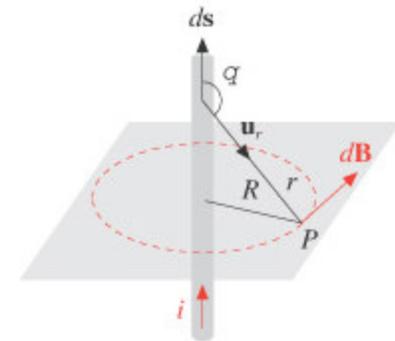
$$\vec{B} = \frac{\mu_0}{4\pi} \frac{q\vec{v} \times \vec{u}_r}{r^2}$$

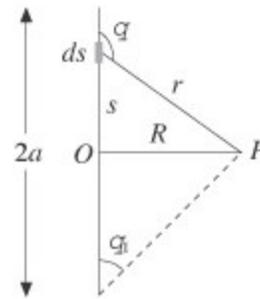
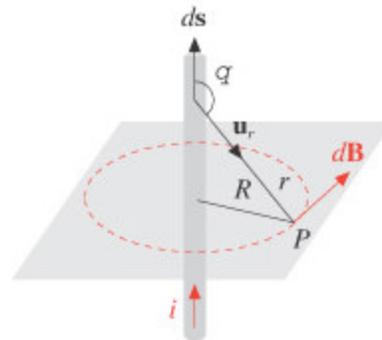
Legge di Biot-Savart

Filo conduttore rettilineo di lunghezza $2a$, corrente i

$$d\vec{B} = \frac{\mu_0}{4\pi} i \frac{d\vec{s} \times \vec{u}_r}{r^2}$$

$$dB = \frac{\mu_0}{4\pi} i \frac{ds \sin\theta}{r^2}$$





$$r \sin(\pi - \theta) = r \sin \theta = R \Rightarrow \frac{1}{r^2} = \frac{\sin^2 \theta}{R^2}$$

$$s \operatorname{tg}(\pi - \theta) = -s \operatorname{tg} \theta = R \Rightarrow ds = \frac{R d\theta}{\sin^2 \theta}$$

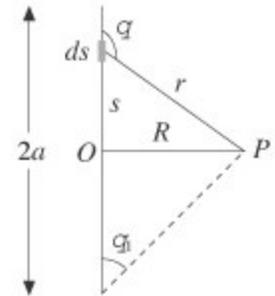
$$\begin{aligned} dB &= \frac{\mu_0}{4\pi} i \frac{ds \sin\theta}{r^2} = \\ &= \frac{\mu_0}{4\pi} i \frac{\sin\theta d\theta}{R} = \\ &= -\frac{\mu_0}{4\pi} i \frac{d(\cos\theta)}{R} \end{aligned}$$

Tratto di filo di lunghezza a produce il campo di modulo:

$$B_a = -\frac{\mu_0 i}{4\pi R} \int_{\cos \theta_1}^{\cos \frac{\pi}{2}} d(\cos \theta) = \frac{\mu_0 i \cos \theta_1}{4\pi R}$$

$$B = 2B_a = \frac{\mu_0 i a}{2\pi R \sqrt{R^2 + a^2}}$$

$$\vec{B} = \frac{\mu_0 i a}{2\pi R \sqrt{R^2 + a^2}} \vec{u}_\phi$$



Legge di Biot-Savart

$$\vec{B} = \frac{\mu_0 i a}{2\pi R \sqrt{R^2 + a^2}} \vec{u}_\phi \xrightarrow{a \rightarrow \infty} = \frac{\mu_0 i}{2\pi R} \vec{u}_\phi$$

$$\vec{B} = \frac{\mu_0 i}{2\pi R} \vec{u}_\phi$$

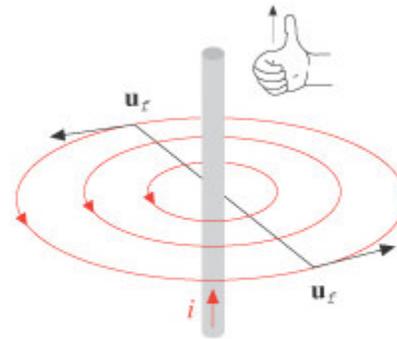
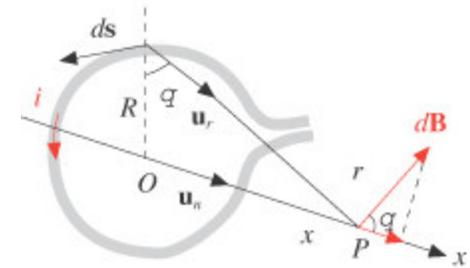


Figura 7.4

Spira circolare

$$dB = \frac{\mu_0}{4\pi} i \frac{ds \operatorname{sen} \theta}{r^2} = \frac{\mu_0}{4\pi} i \frac{ds}{r^2}$$



$$dB_x = \frac{\mu_0}{4\pi} i \frac{ds \cos \theta}{r^2}$$

$$\vec{B} = \oint \frac{\mu_0}{4\pi} i \frac{\cos \theta}{r^2} ds \vec{u}_n = \frac{\mu_0}{4\pi} i \frac{\cos \theta}{r^2} 2\pi R \vec{u}_n$$

posto

$$r^2 = x^2 + R^2$$

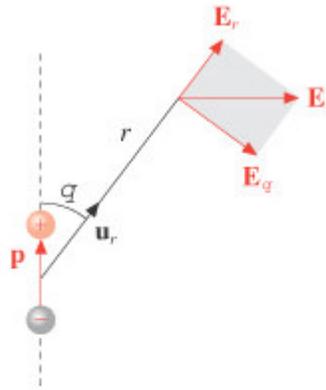
$$\cos \theta = \frac{R}{r}$$

$$\vec{B} = \frac{\mu_0 i R^2}{2(x^2 + R^2)^{3/2}} \vec{u}_n$$

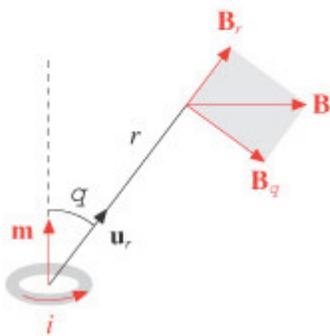
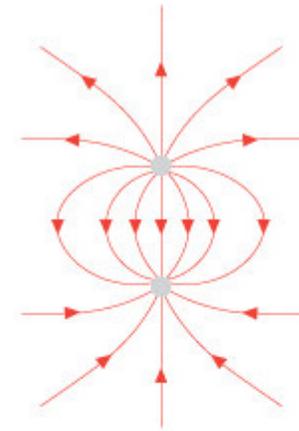
Per $x \gg R$

$$\vec{B} = \frac{\mu_0 i R^2}{2(x^2 + R^2)^{3/2}} \vec{u}_n \rightarrow \frac{\mu_0 i R^2}{2x^3} \vec{u}_n = \frac{\mu_0}{4\pi} \frac{2\vec{m}}{r^3}$$

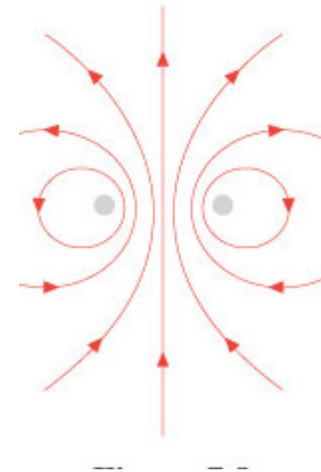
$$\vec{m} = i\Sigma \vec{u}_n = i\pi R^2 \vec{u}_n$$



$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{2\vec{p}}{r^3}$$



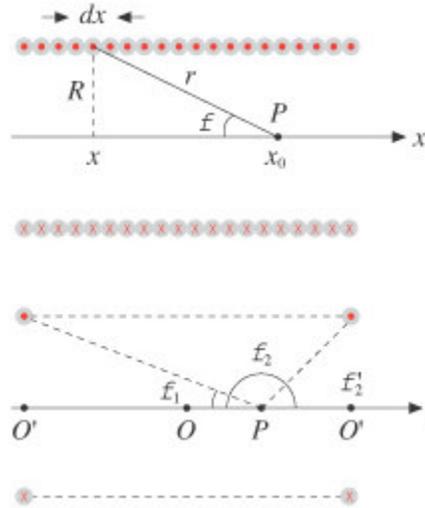
$$\vec{B} = \frac{\mu_0}{4\pi} \frac{2\vec{m}}{r^3}$$



$$\vec{B} = \frac{\mu_0}{4\pi} \frac{m}{r^3} (2\cos\theta \vec{u}_r + \sin\theta \vec{u}_\theta)$$

Solenoide rettilineo

$$dB = \frac{\mu_0 i R^2 n}{2r^3} dx$$



$$r \sin \phi = R$$

$$x - x_0 = -R \cot \phi$$

$$dx = \frac{R d\phi}{\sin^2 \phi}$$

$$dB = \frac{\mu_0 i n}{2} \sin \phi d\phi$$

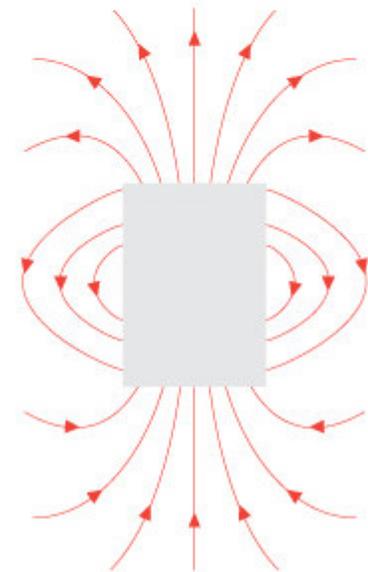
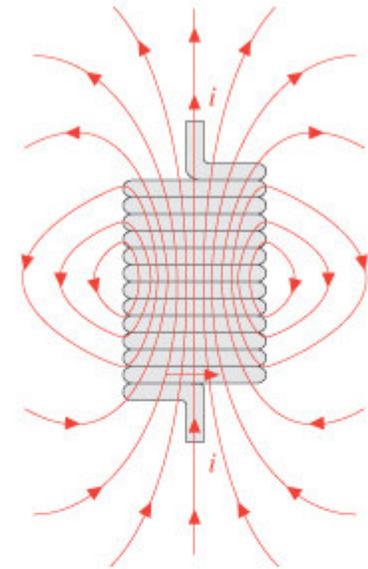
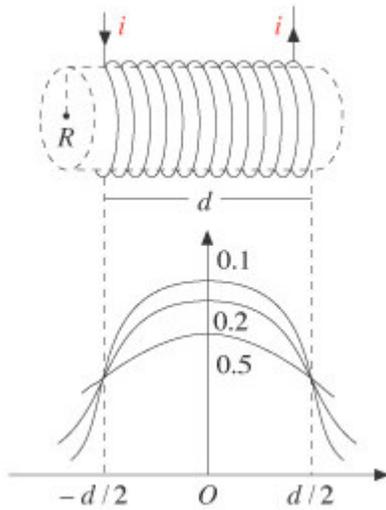
$$\begin{aligned}
 B &= \int_{\phi_1}^{\phi_2} dB = \frac{\mu_0 in}{2} \int_{\phi_1}^{\phi_2} \text{sen } \phi d\phi = \\
 &= \frac{\mu_0 in}{2} (\cos \phi_1 + \cos \phi_2')
 \end{aligned}$$

Misurando rispetto al centro del solenoide ($x=0$)

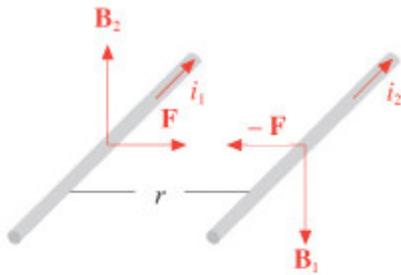
$$B = \frac{\mu_0 in}{2} \left(\frac{d + 2x}{\sqrt{(d + 2x)^2 + 4R^2}} + \frac{d - 2x}{\sqrt{(d - 2x)^2 + 4R^2}} \right)$$

Per $d \gg R$

$$B_{\infty} = \mu_0 i n$$

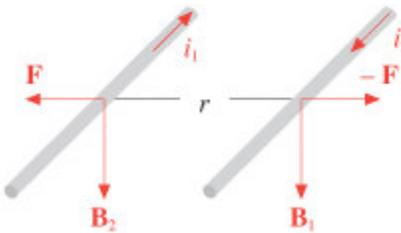


Definizione di Ampère



$$d\vec{F}_{12} = i_2 d\vec{s} \times \vec{B}_1 = i_2 ds \vec{u}_2 \times \vec{B}_1$$

$$\vec{B}_1 = \frac{\mu_0 i_1}{2\pi r} \vec{u}_\phi$$



$$\frac{dF_{12}}{ds} = F_{12} = F_{21} = \frac{\mu_0 i_1 i_2}{2\pi r}$$

Definizione di Ampère

Ha l'intensità di 1A quella corrente che circolando in due fili rettilinei paralleli distanti $r = 1\text{ m}$ dà luogo ad una forza $F = \mu/2\pi = 2 \cdot 10^{-7}\text{ N}$ per metro di ciascun conduttore.

Relazione circuitazionale di Ampère

Le proprietà fondamentali di un campo vettoriale consistono nell'essere o no solenoidale, nell'essere o no conservativo.

Relazione circuitazionale di Ampère

Si dice che un campo vettoriale è solenoidale quando il flusso di quel vettore uscente da una qualunque superficie chiusa è nullo

Relazione circuitazionale di Ampère

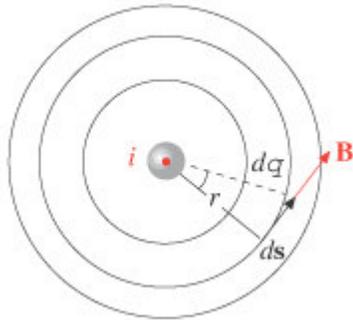
Si dice che un campo vettoriale è conservativo quando la circuitazione di quel vettore lungo una qualunque linea chiusa è nulla

Relazione circuitazionale di Ampère

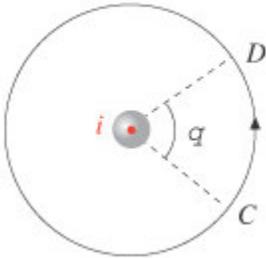
Il vettore \mathbf{B} è o non è conservativo???

Relazione circuitazionale di Ampère

$$\vec{B} \bullet d\vec{s} = \frac{\mu_0 i}{2\pi r} ds = \frac{\mu_0 i}{2\pi} d\theta$$

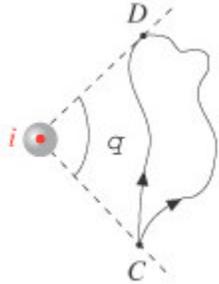


$$\int_C^D \vec{B} \bullet d\vec{s} = \frac{\mu_0 i}{2\pi} \int_C^D d\theta = \frac{\mu_0 i}{2\pi} \theta$$

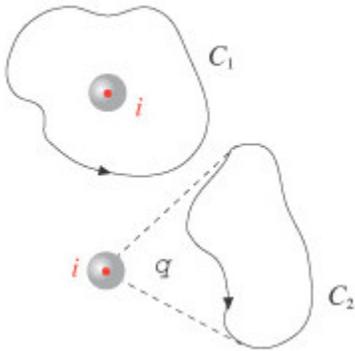


$$\int_D^C \vec{B} \bullet d\vec{s} = -\frac{\mu_0 i}{2\pi} \theta$$

Relazione circuitazionale di Ampère



$$\oint \vec{B} \cdot d\vec{s} = \frac{\mu_0 i}{2\pi} \oint d\theta$$



$$\oint \vec{B} \cdot d\vec{s} = \frac{\mu_0 i}{2\pi} 2\pi = \mu_0 i$$

$$\oint \vec{B} \cdot d\vec{s} = 0$$

Relazione circuitazionale di Ampère

o

Legge di Ampère

La circuitazione di \mathbf{B} lungo una linea chiusa qualunque è uguale alla somma delle correnti concatenate, moltiplicata per μ_0

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i$$

Per il teorema di Stokes

$$\oint \vec{B} \cdot d\vec{s} = \int_{\Sigma} \vec{\nabla} \times \vec{B} \cdot \vec{u}_n d\Sigma$$

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i =$$

$$= \mu_0 \int_{\Sigma} \vec{j} \cdot \vec{u}_n d\Sigma = \int_{\Sigma} \vec{\nabla} \times \vec{B} \cdot \vec{u}_n d\Sigma$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j}$$

Legge di Ampère

globale

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i$$

locale

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j}$$

Forma globale

$$\oint_{\Sigma} \vec{E} \cdot \vec{u}_n d\Sigma = \frac{q}{\epsilon_0}$$

$$\oint \vec{E} \cdot d\vec{S} = 0$$

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i$$

Forma locale

$$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho$$

$$\vec{\nabla} \times \vec{E} = 0$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j}$$

Legge di Gauss

$$\oint_{\Sigma} \vec{B} \cdot \vec{u}_n d\Sigma = 0$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

B è solenoidale

Forma globale

$$\oint_{\Sigma} \vec{E} \cdot \vec{u}_n d\Sigma = \frac{q}{\epsilon_0}$$

$$\oint \vec{E} \cdot d\vec{S} = 0$$

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 i$$

$$\oint_{\Sigma} \vec{B} \cdot \vec{u}_n d\Sigma = 0$$

Forma locale

$$\vec{\nabla} \cdot \vec{E} = \frac{1}{\epsilon_0} \rho$$

$$\vec{\nabla} \times \vec{E} = 0$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{j}$$

$$\vec{\nabla} \cdot \vec{B} = 0$$

Equazioni di Maxwell

$$\oint \vec{E} \cdot \vec{u}_n d\Sigma = \frac{q}{\epsilon_0}$$

$$\oint \vec{B} \cdot \vec{u}_n d\Sigma = 0$$

$$\oint \vec{E} \cdot d\vec{s} = -\frac{d\Phi(\vec{B})}{dt}$$

$$\oint \vec{B} \cdot d\vec{s} = \mu_0 \left(i + \epsilon_0 \frac{d\Phi(\vec{E})}{dt} \right)$$